

Physics with extra dimensions

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Part 2

- Large extra dimensions and low scale gravity
- Experimental signals
- Warped extra dimensions

Low scale gravity

Extra large $\perp$ dimensions can explain the apparent weakness of gravity

total force = observed force $\times$ volume $\perp$

total force $\simeq \mathcal{O}(1)$ at 1 TeV n dimensions of size $R_\perp$

$\Rightarrow$ volume $R_\perp^n = 10^{32}$ in TeV^{-1} units

$n = 1 : R_\perp \simeq 10^8 \text{ km}$ excluded

$n = 2 : R_\perp \simeq 0.1 \text{ mm}$ (10^{-12} GeV) possible

$n = 6 : R_\perp \simeq 10^{-13} \text{ mm}$ (10^{-2} GeV)

- distances $> R_\perp$: gravity 3d

however for $< R_\perp$: gravity $(3+n)$ d [4]

- strong gravity at $10^{-16} \text{ cm} \leftrightarrow 10^3 \text{ GeV}$

10^{30} times stronger than thought previously! [7]

Low scale gravity

Extra large $\perp$ dimensions can explain the apparent weakness of gravity

total force = observed force $\times$ volume $\perp$

$$\begin{array}{ccccc} \uparrow & & \uparrow & & \uparrow \\ G_N^* E^{2+n} & = & G_N E^2 & \times & V_{\perp} E^n \end{array}$$

$G_N^* = M_*^{-(2+n)}$: $(4+n)$ -dim gravitational constant

total force $\simeq \mathcal{O}(1)$ at 1 TeV

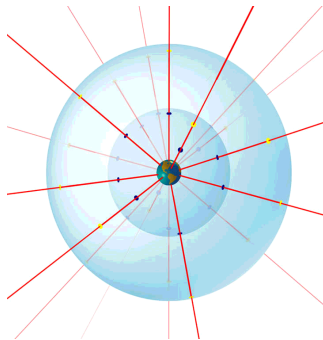
n dimensions of size $R_{\perp}$

$$\Rightarrow V_{\perp} = R_{\perp}^n$$

$$\Rightarrow M_P^2 = M_*^{2+n} R_{\perp}^n \text{ for } M_* \simeq 1 \text{ TeV} \Rightarrow (R_{\perp} M_*)^n \sim 10^{32}$$

Gravity modification at submillimeter distances

Newton's law: force decreases with area



3d: force $\sim 1/r^2$

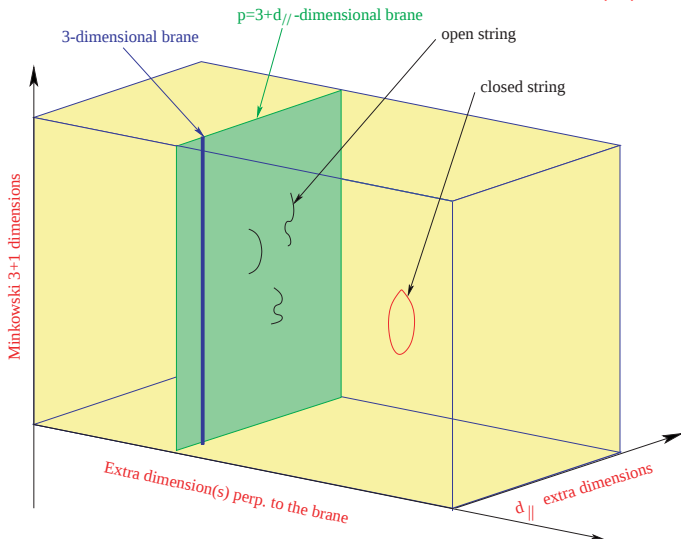
(3+n)d: force $\sim 1/r^{2+n}$

observable for $n = 2$: $1/r^4$ with $r \ll .1$ mm [2]

Braneworld

2 types of compact extra dimensions:

- parallel ($d_{\parallel}$): $\lesssim 10^{-16}$ cm (TeV) [6]
- transverse ($\perp$): $\lesssim 0.1$ mm (meV)



Experimental predictions

- No little hierarchy problem:

radiative electroweak symmetry breaking with no logs

$$\Lambda \sim \text{a few TeV} \text{ and } m_H^2 = \text{a loop factor} \times \Lambda^2$$

- particle accelerators

- Large TeV dimensions seen by gauge interactions [17]
- Extra large hidden dimensions transverse $\Rightarrow$ strong gravity
- other accelerator signatures

- microgravity experiments

- gravity modifications at short distances
- new submillimeter forces

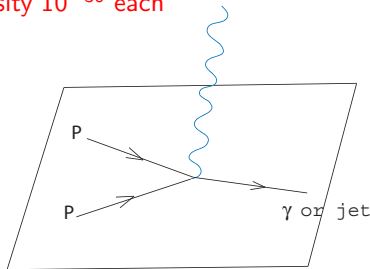
Hidden submillimeter dimensions

⇒ strong gravity at the TeV: gravitational radiation in the bulk

3d: Kaluza Klein gravitons very light

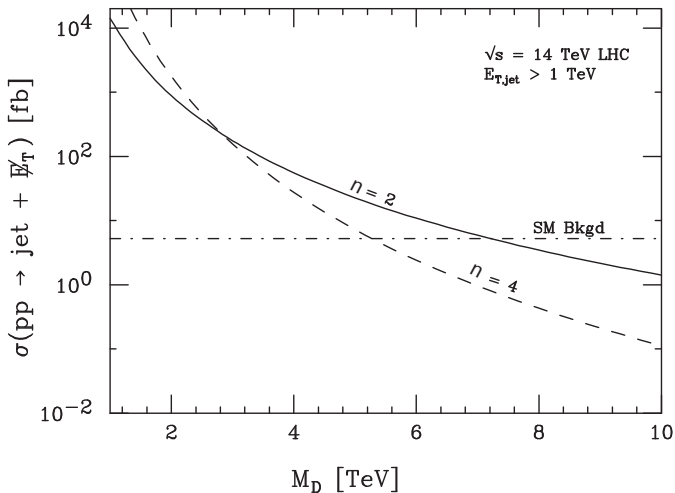
⇒ high energy: huge number of particles produced

LHC: 10^{30} massive gravitons of intensity 10^{-30} each



Signal: missing energy

Angular distribution ⇒ spin of the graviton



- no observation $\Rightarrow R_{\perp} \lesssim 10^{-2} - 10^{-12} \text{ mm}$ ($n = 2 - 6$); 95% CL
- more dimensions $\Rightarrow$ weaker limits

Angular distribution for graviton emission

Sum over KK modes:

$$\begin{aligned}\sum_{k_\perp} &= R^n \int d^n m = \frac{1}{2} \Omega_n R^n \int (m^2)^{(n-2)/2} d m^2 \\ &= \Omega_n \frac{M_p^2}{2 M_*^{(n+2)}} \int (m^2)^{(n-2)/2} d m^2\end{aligned}$$

$$\Omega_n = \frac{2\pi^{n/2}}{\Gamma(n/2)} : \text{volume of } n\text{-sphere of unit radius}$$

Angular distribution for graviton emission

$$\begin{aligned}
 |\mathcal{M}(q\bar{q} \rightarrow gG_m)|^2 &= \frac{1}{18M_p^2} \frac{\alpha_s}{1 - m^2/s} \left[\left(2 - \frac{4ut}{(s - m^2)^2}\right) \left(1 + \left(\frac{m^2}{s}\right)^4\right) \right. \\
 &\quad + \left(2\frac{(s - m^2)^2}{4ut} - 5 + 4\frac{4ut}{(s - m^2)^2}\right) \frac{m^2}{s} \left(1 + \left(\frac{m^2}{s}\right)^2\right) \\
 &\quad \left. + 6\left(\frac{u - t}{s - m^2}\right)^2 \left(\frac{m^2}{s}\right)^2 \right]
 \end{aligned}$$

$$\begin{aligned}
 |\mathcal{M}(qg \rightarrow qG_m)|^2 &= \frac{\alpha_s}{48M_p^2} \frac{-t/s(1 - m^2/s)}{(1 - m^2/t)^2} \left[\left(2 - \frac{4us}{(t - m^2)^2}\right) \left(1 + \left(\frac{m^2}{t}\right)^4\right) \right. \\
 &\quad + \left(2\frac{(t - m^2)^2}{4us} - 5 + 4\frac{4us}{(t - m^2)^2}\right) \frac{m^2}{t} \left(1 + \left(\frac{m^2}{t}\right)^2\right) \\
 &\quad \left. + 6\left(\frac{s - u}{t - m^2}\right)^2 \left(\frac{m^2}{t}\right)^2 \right]
 \end{aligned}$$

$$t, u = -\frac{1}{2}s(1 - m^2/s)(1 \mp \cos \theta)$$

Cross sections

$$\left. \begin{array}{l} |\mathcal{M}(gg \rightarrow gg)|^2 \quad , \quad |\mathcal{M}(gg \rightarrow q\bar{q})|^2 \\ |\mathcal{M}(q\bar{q} \rightarrow gg)|^2 \quad , \quad |\mathcal{M}(qg \rightarrow qg)|^2 \end{array} \right\} \quad \begin{array}{l} \text{model independent} \\ \text{for any compactification} \end{array}$$

$$|\mathcal{M}(gg \rightarrow gg)|^2 = g_{YM}^4 \left(\frac{1}{s^2} + \frac{1}{t^2} + \frac{1}{u^2} \right) \times \left[\frac{9}{4} (s^2 V_s^2 + t^2 V_t^2 + u^2 V_u^2) - \frac{1}{3} (sV_s + tV_t + uV_u)^2 \right]$$

$$|\mathcal{M}(gg \rightarrow q\bar{q})|^2 = g_{YM}^4 \frac{t^2 + u^2}{s^2} \left[\frac{1}{6} \frac{1}{tu} (tV_t + uV_u)^2 - \frac{3}{8} V_t V_u \right] \quad M_s = 1$$

$$V_s = -\frac{tu}{s} \quad B(t, u) = 1 - \frac{2}{3}\pi^2 tu + \dots \quad V_t : s \leftrightarrow t \quad V_u : s \leftrightarrow u$$

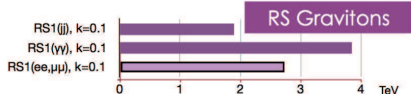
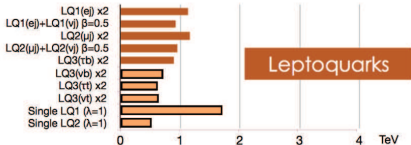
YM limits agree with e.g. book "*Collider Physics*" by Barger, Phillips

Limits on $R_{\perp}$ in mm

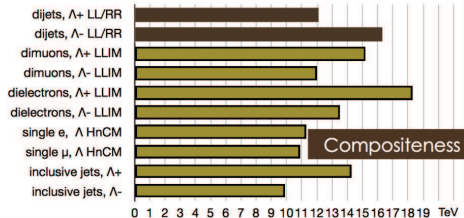
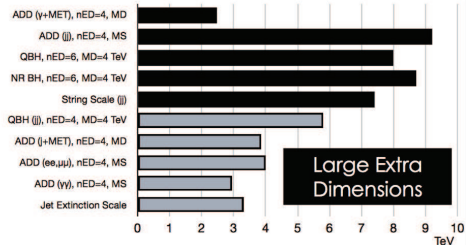
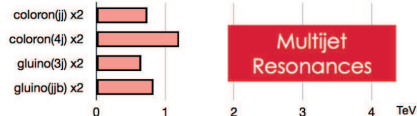
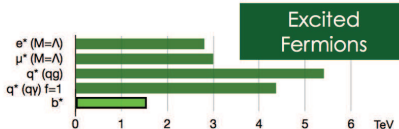
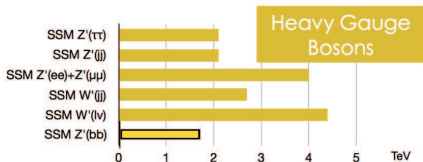
Experiment	$R_{\perp}(n=2)$	$R_{\perp}(n=4)$	$R_{\perp}(n=6)$
Collider bounds			
LEP 2	4.8×10^{-1}	1.9×10^{-8}	6.8×10^{-11}
Tevatron	5.5×10^{-1}	1.4×10^{-8}	4.1×10^{-11}
LHC	4.5×10^{-3}	5.6×10^{-10}	2.7×10^{-12}
Astrophysics/cosmology bounds			
SN1987A	3×10^{-4}	1×10^{-8}	6×10^{-10}
COMPTEL	5×10^{-5}	-	-

present LHC bounds: $M_* \gtrsim 3 - 5$ TeV

13 TeV 8 TeV



CMS Preliminary



Supernova constraints

cooling due to graviton production e.g. $NN \rightarrow NN + \text{graviton}$

number of gravitons: $\sim (TR_{\perp})^n$

$$\Rightarrow \text{production rate: } P_{\text{gr}} \sim \frac{1}{M_P^2} (TR_{\perp})^n \sim \frac{T^n}{M_*^{(2+n)}}$$

$$P_{\text{gr}} < P_\nu \Rightarrow M_*|_{n=2} \gtrsim 50 \text{ TeV} \Rightarrow M_s \gtrsim 10 \text{ TeV}$$

R-neutrinos: in the bulk

R-neutrino: $\nu_R(x, y)$ y : bulk coordinates

$$S_{int} = g_s \int d^4x H(x) L(x) \nu_R(x, y=0)$$

$\langle H \rangle = v \Rightarrow$ mass-term: $\frac{g_s v}{R_\perp^{n/2}} \nu_L \nu_R^0 \leftarrow 4d \text{ zero-mode}$

Dirac neutrino masses: $m_\nu \simeq \frac{g_s v}{R_\perp^{n/2}} \simeq v \frac{M_*}{M_p}$

$$\simeq 10^{-3} - 10^{-2} \text{ eV for } M_* \simeq 1 - 10 \text{ TeV}$$

$m_\nu \ll 1/R_\perp \Rightarrow$ KK modes unaffected

Black hole production

String-size black hole energy threshold : $M_{\text{BH}} \simeq M_s/g_s^2$

- string size black hole: $r_H \sim l_s = M_s^{-1}$
- black hole mass: $M_{\text{BH}} \sim r_H^{d-3}/G_N$ $G_N \sim l_s^{d-2} g_s^2$

weakly coupled theory $\Rightarrow$ strong gravity effects occur much above M_s , M_*

$$g_s \sim 0.1 \text{ (gauge coupling)} \quad \Rightarrow \quad M_{\text{BH}} \sim 100 M_s$$

Comparison with Regge excitations : $M_n = M_s \sqrt{n} \Rightarrow$

production of $n \sim 1/g_s^4 \sim 10^4$ string states before reach M_{BH} [6]

Large TeV dimensions

longitudinal dimensions: $R^{-1} \lesssim M_s \Rightarrow R^{-1}$ first scale of new physics

increasing the energy

- could happen for some of the internal dims
- explain coupling constant ratios g_2/g_3
- susy breaking
- fermion masses displace light generations

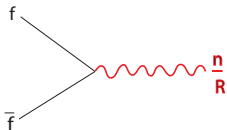
Massive tower of Kaluza Klein modes for Standard Model particles

$$M_n^2 = M_0^2 + \frac{n^2}{R^2} \quad ; \quad n = \pm 1, \pm 2, \dots$$

$\Rightarrow$ excited states of photon, $W^\pm$, Z, gluons

Localized fermions (on 3-brane intersections)

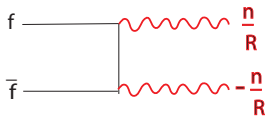
⇒ single production of KK modes



- strong bounds indirect effects: $R^{-1} \gtrsim 4 \text{ TeV}$
- new resonances but at most $n = 1$

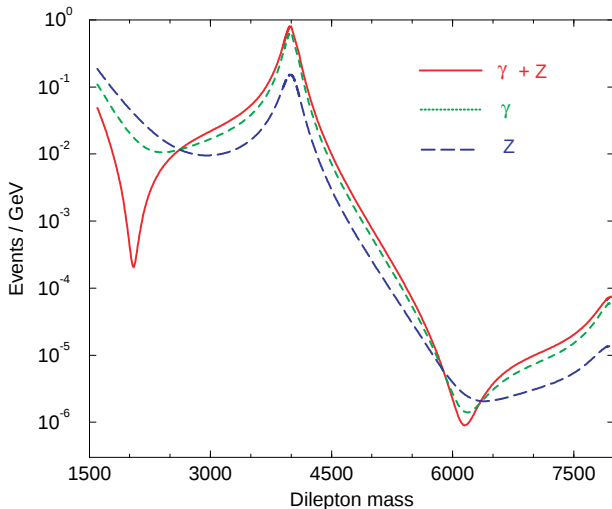
Otherwise KK momentum conservation [21]

⇒ pair production of KK modes (universal dims)



- weak bounds $R^{-1} \gtrsim 500 \text{ GeV}$
- no resonances
- lightest KK stable ⇒ dark matter candidate

$$R^{-1} = 4 \text{ TeV}$$



- no observation in dijets $\Rightarrow R^{-1} \gtrsim 20 \text{ TeV}$; 95% CL
- more than one dimension $\Rightarrow$ stronger limits

KK W -production at LHC in the $l\nu$ channel [18]

